

Study of Color Transparency in Exclusive Vector Meson Electroproduction off Nuclei

Hall C Summer Workshop

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COLLABORATION

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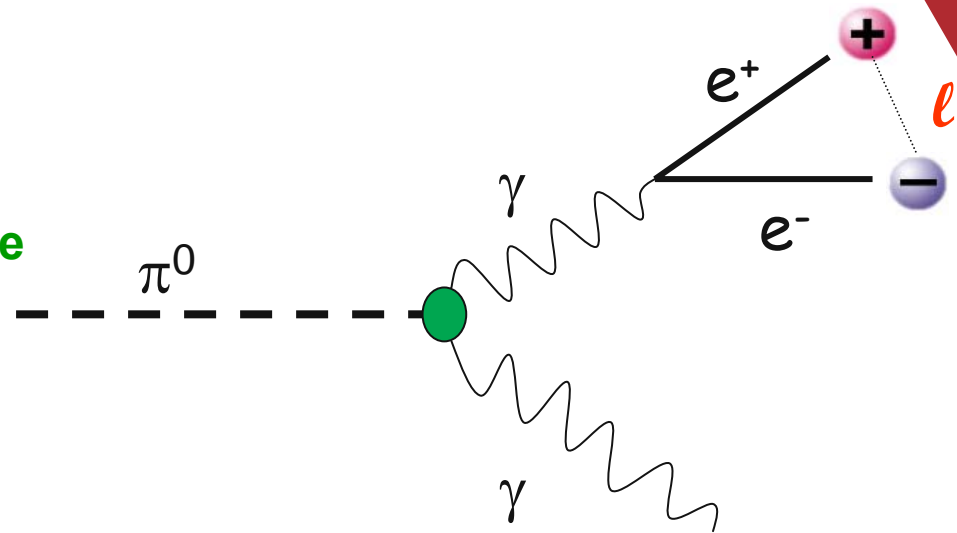
Overview

- **Physics Motivation**
- **EG2 Experiment**
- **Preliminary Results**
- **Summary and Outlook**



Origin of CT

- Discovery by Perkins (1955) of the (Dalitz) decays in emulsion of π^0 (~ 200 GeV) produced in cosmic rays $\pi^0 \rightarrow e^+ e^- \gamma$



- The ionization produced by the pair was small near the decay point, increasing with distance from vertex
- This surprising observation was quickly interpreted by Chudakov (1955) in the framework of QED: A pair of oppositely charged particles interacts in the medium with a dipole cross-section

\Rightarrow this cross-section ($\sigma \cong \ell^2$) vanishes near the creation point

In early 80's, **Brodsky and Mueller** applied the notion of transparency to **QCD** and to **color charge**

Color Transparency is a spectacular **prediction of QCD**:
Under the right conditions, the nuclear matter will allow the transmission of hadrons with reduced attenuation. Such a phenomenon is totally **unexpected in a hadronic picture of strongly interacting matter**, but **straightforward in quark**

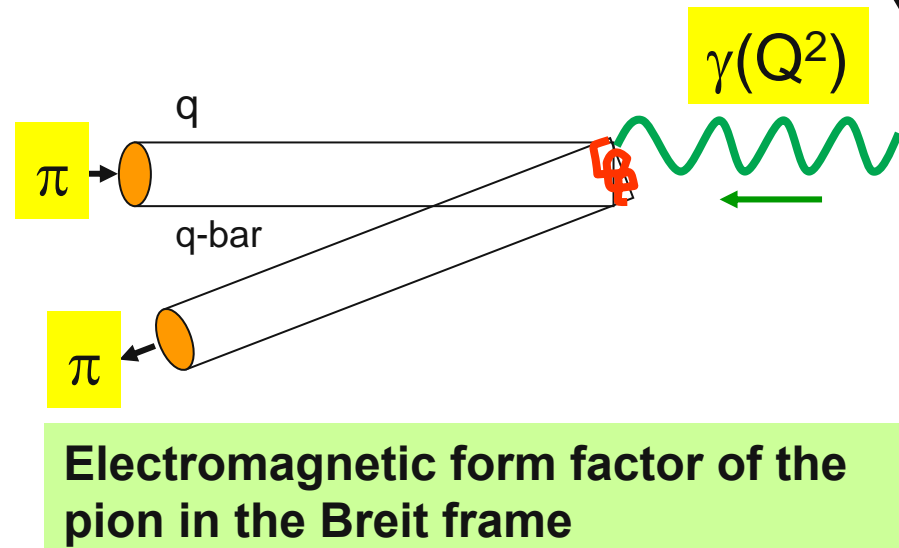
gluon basis, this is one of the features which makes it so

Right Conditions \Rightarrow Selection of Point Like Configurations (PLC) via hard exclusive processes

Hard: high momentum transfer

Exclusive: completely determined initial and final states.

Elastic processes are special cases



◆ Unless the struck quark shares the momentum transfer with the other quark, the pion fragments and the reaction is inelastic

◆ As Q increases, the exchange of the gluon has to be fast.

Causality (no interaction is faster than speed of light)

\Rightarrow the quark's pair has to be localized within a transverse size of $1/Q$

Color screening: PLC experiences reduced interaction in the nucleus

In QCD the color field of a color neutral object vanishes as the size of the object is reduced

Because

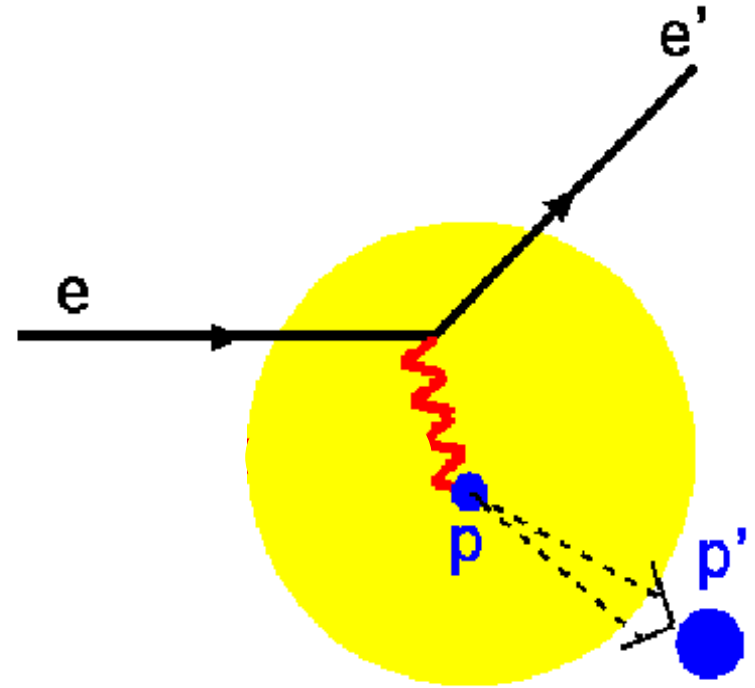
The field of individual quarks and gluon cancel each other as the size is reduced by analogy to QED

Therefore

The interaction cross-section has a dipole form

$$\sigma \cong \ell^2$$

ℓ is the separation between the constituents



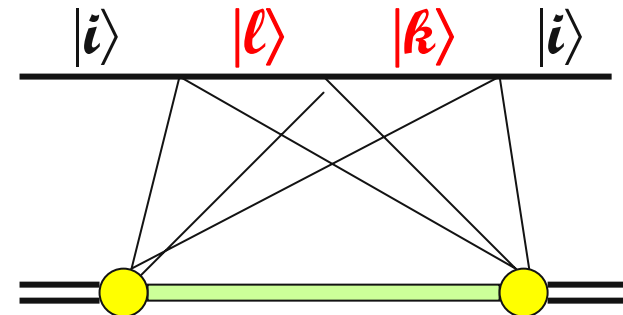
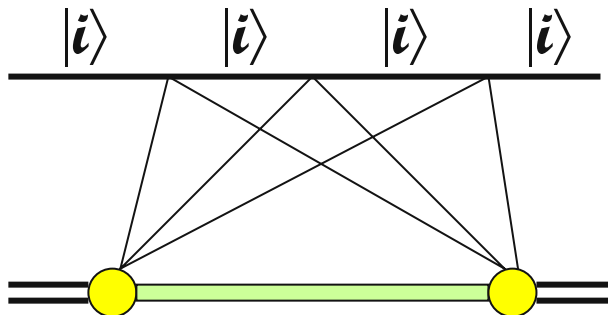
Hadronic point of view of CT & PLC formation time

- The point like aspect is seen as a consequence of coherent superposition of large number of resonances with specific weights
- CT is understood as a coherence of the scattering of these resonances inside the nucleus
- The formation of PLC is a function of the typical excitation energy of the system

$$\frac{1}{\tau_f} \cong \sqrt{E_h^2 + M^{*2}} - \sqrt{E_h^2 + M^2}$$

Assuming $v \cong E_h \cong p$

$$\tau_f \cong \frac{2v}{M^{*2} - M^2}$$



What can we learn from studying CT ?

$$|\text{Meson}\rangle = Z_0 |q \bar{q}\rangle + Z_1 |q \bar{q} q \bar{q}\rangle +$$

.....

- PLC is by definition a product of short distances: it can only come from valence component (higher order are reduced by a factor α_s)
- CT mechanism selects the simplest component of the hadron wave-function. By analogy to lattice QCD, we are in the “quenched approximation”
- All the physics programs build around CT idea would allow us not only to access special configurations of the hadron wave-function but also to study how this configuration dresses with time to form the asymptotic wave-function of the hadron with all its complexity
- We are here in the heart of the dynamics of confinement !

The nucleus, a unique laboratory of quark dynamics

✿ **Characteristic proper time scale is $\tau_0 \sim 1$ fm**

τ_0 is the **time** needed by a quark to travel distances typical of the confined systems

✿ Taking into account **Lorentz dilation**, the proper time scales in

the Lab frame become $\tau = (E/M) \tau_0 \sim \text{few fm}$

✿ **The only medium available for these scales is the nucleus !**

✿ **The nucleus is playing the role of the bubble chamber !**



ρ^0 electroproduction on nuclei

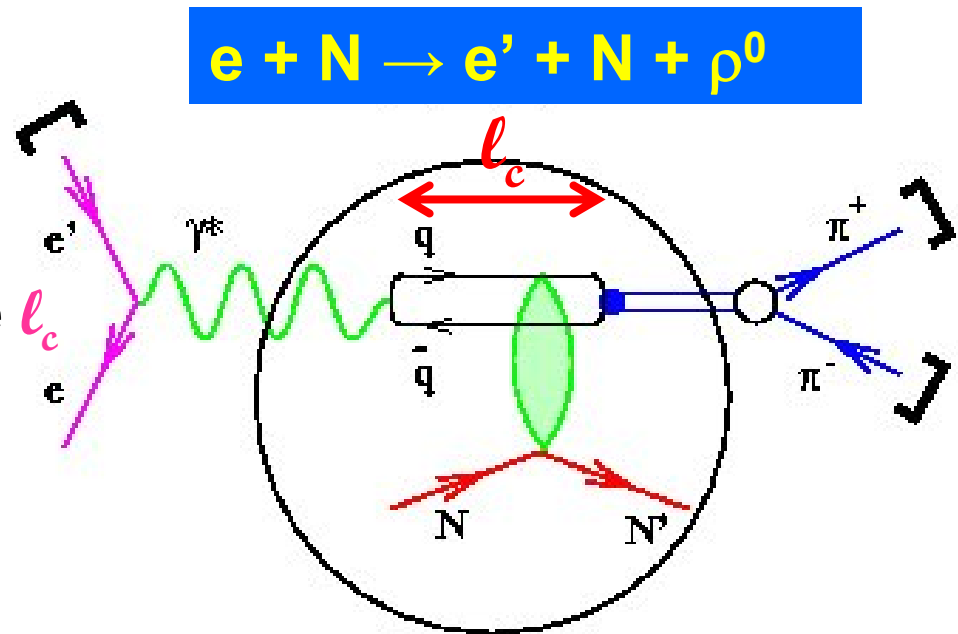
Detected particles are :
scattered electron and the
 π^+ and π^- from ρ^0 decay

Finite propagation distance
(lifetime) of the $(q, q\text{-bar})$
virtual state

$$\ell_c = 2v/(M^2 + Q^2)$$

M is the mass of the vector meson

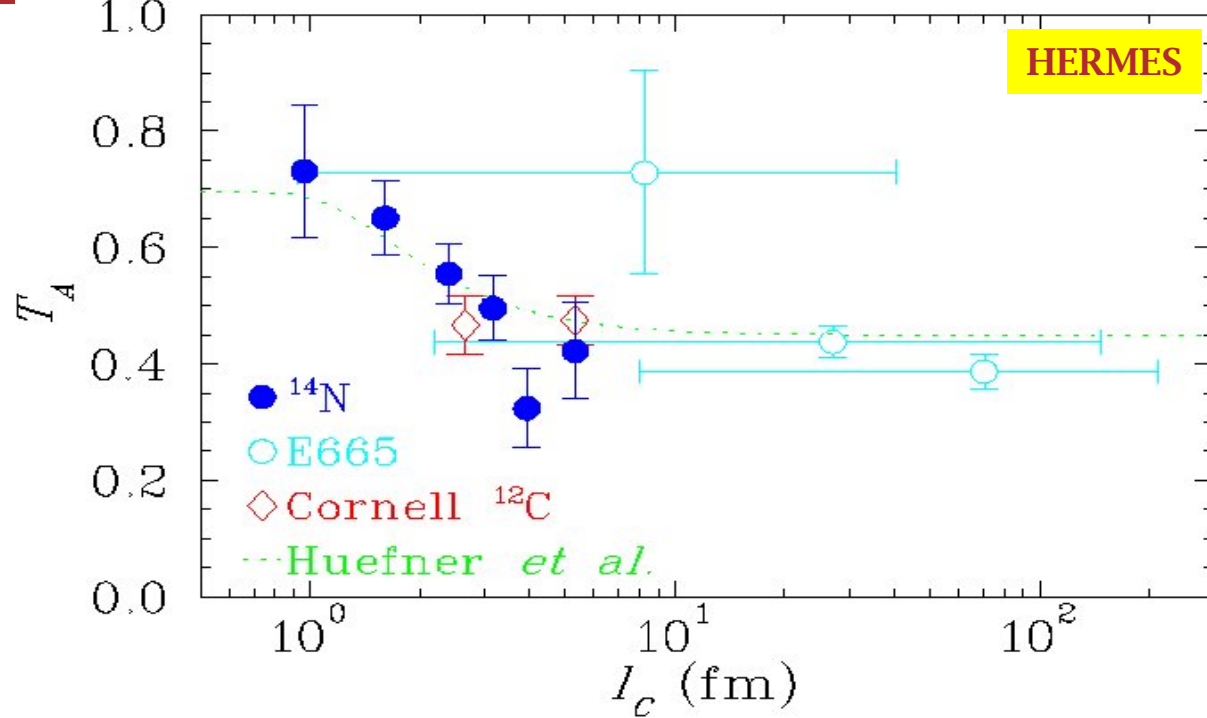
v is the energy transferred by the electron



What Could mimic CT signal ?

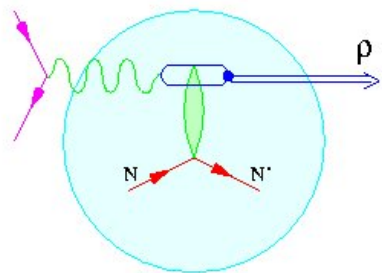
Coherence Length

$$\ell_c = 2v/(M_v^2 + Q^2)$$



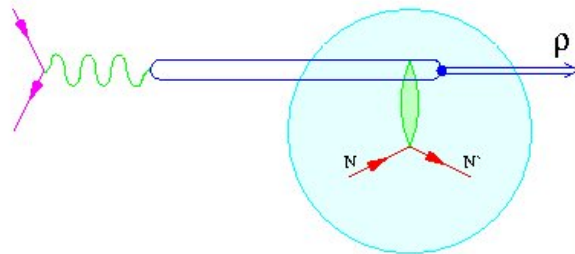
Coherence length effect (CL): Q^2 increases $\Rightarrow T_A$ increases

Electromagnetic ISI



Small ℓ_c

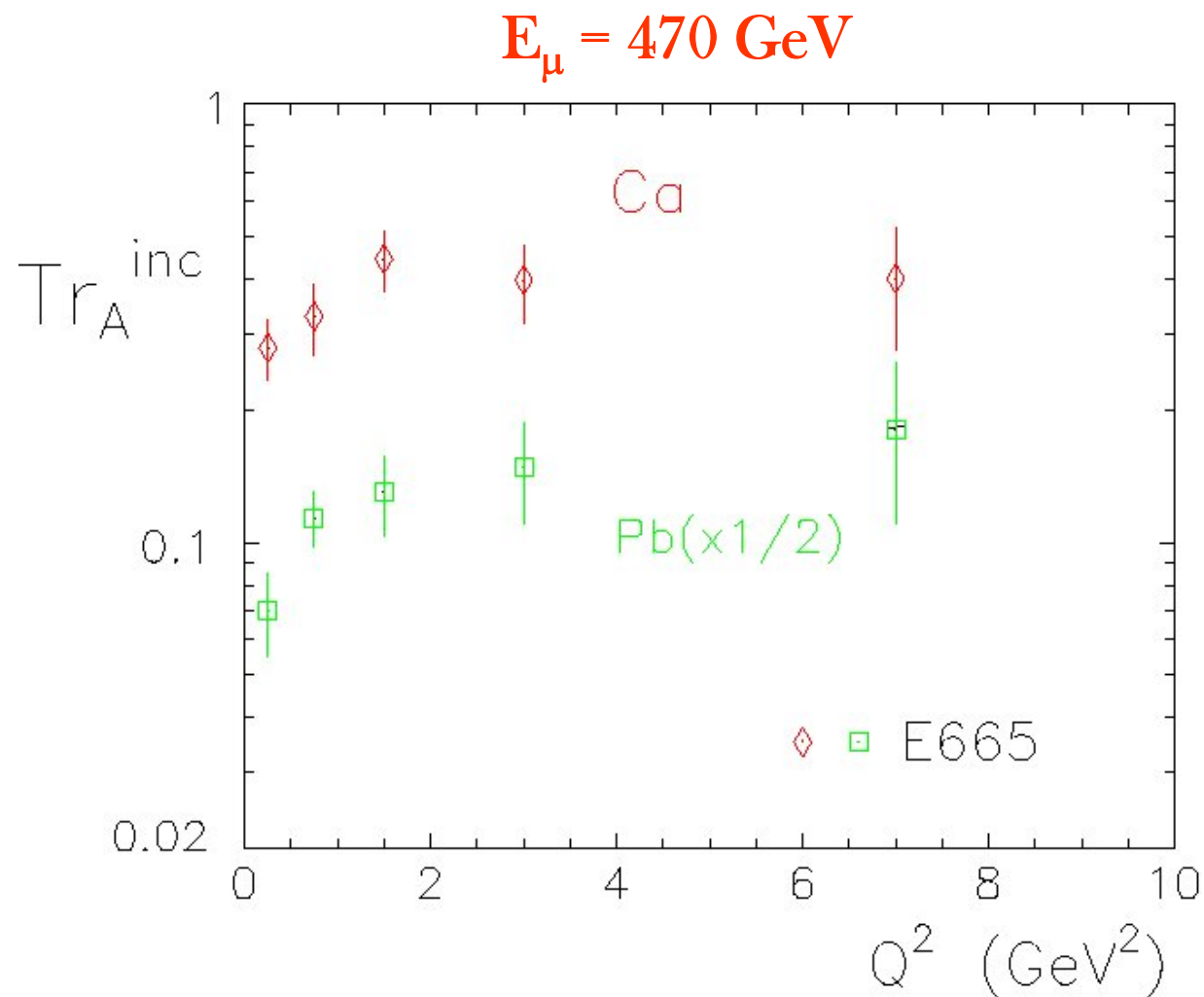
Hadronic ISI



Large ℓ_c

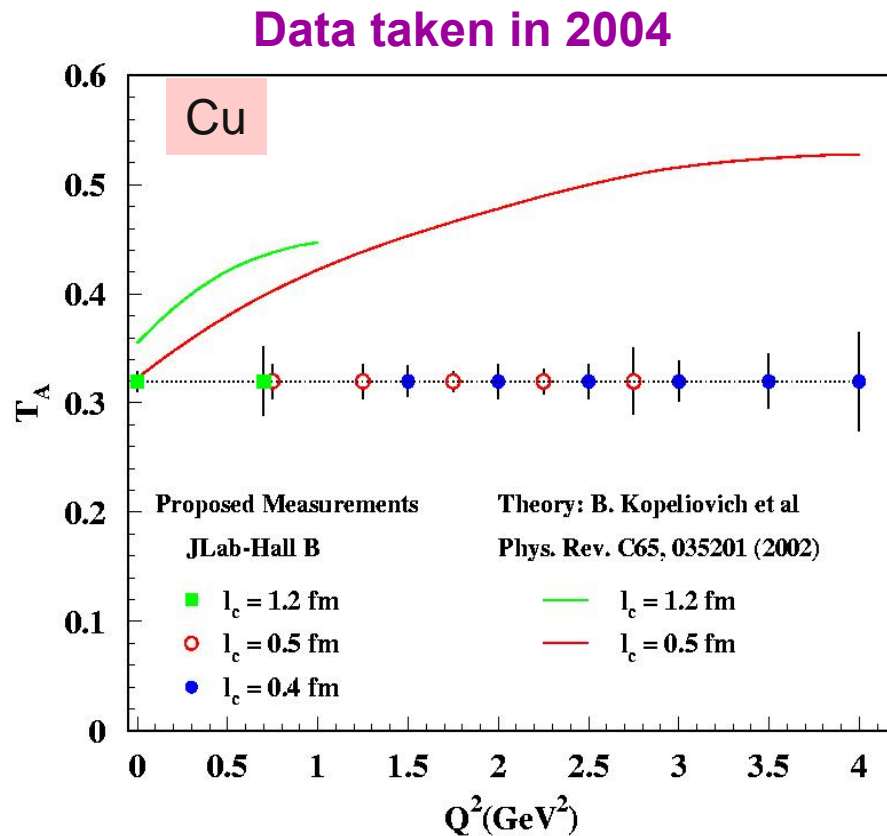
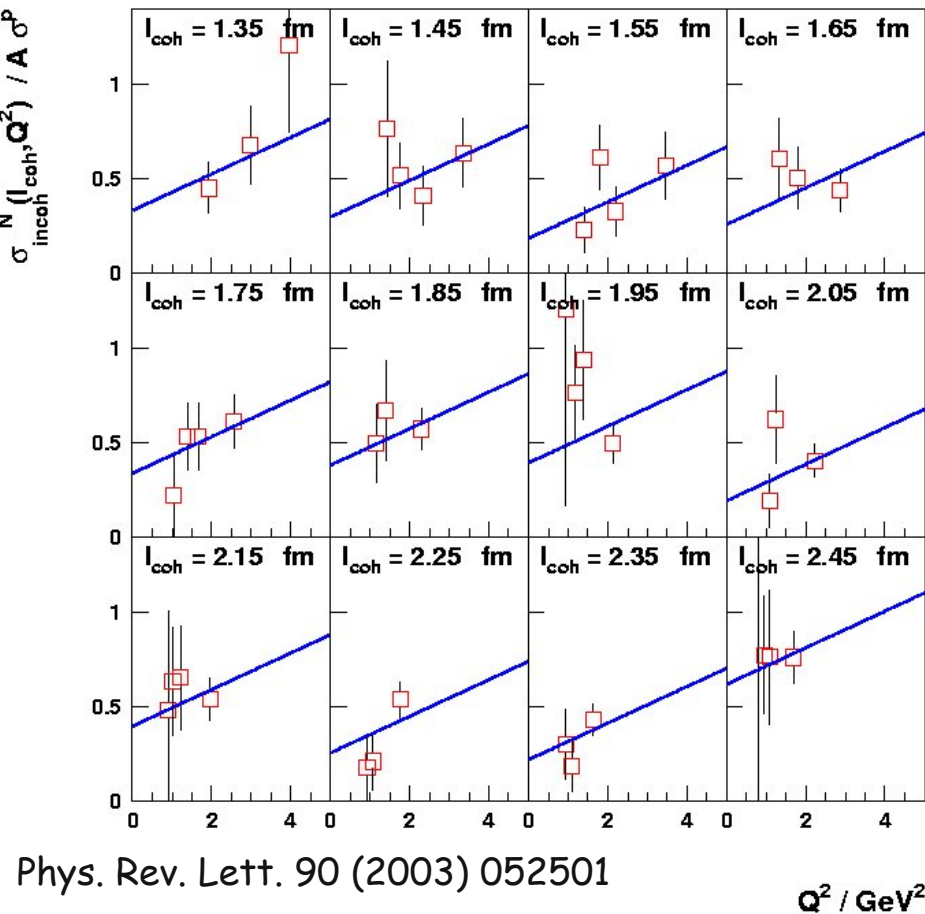
- **Coherence Length effect can mimic CT signal**
- **To be safe, one should keep ℓ_c fixed and measures the Q^2 dependence of T_A**

FNAL E665 experiment



Adams et al. PRL74 (1995) 1525

ρ^0 electroproduction at fixed CL

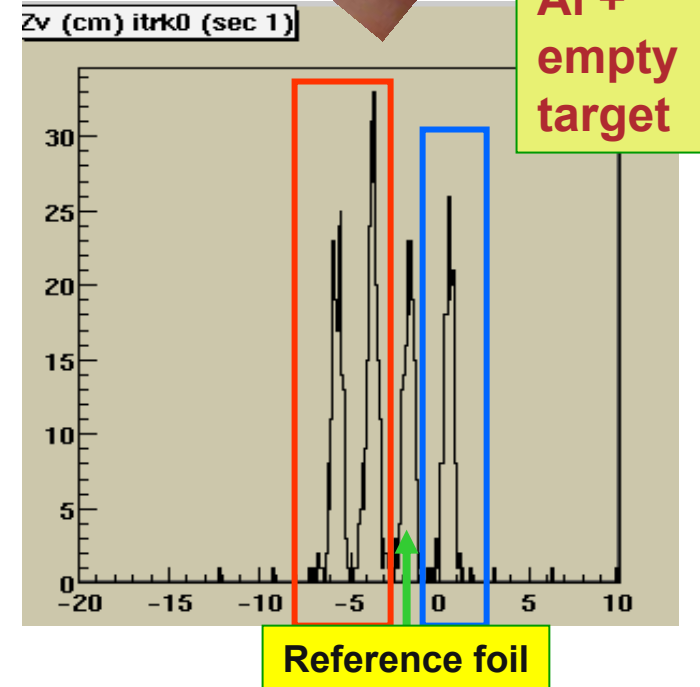
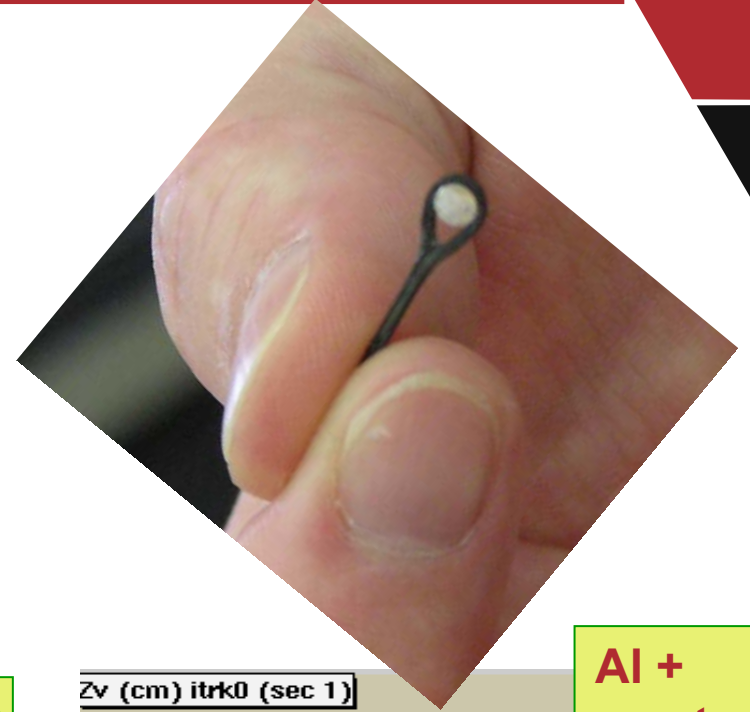
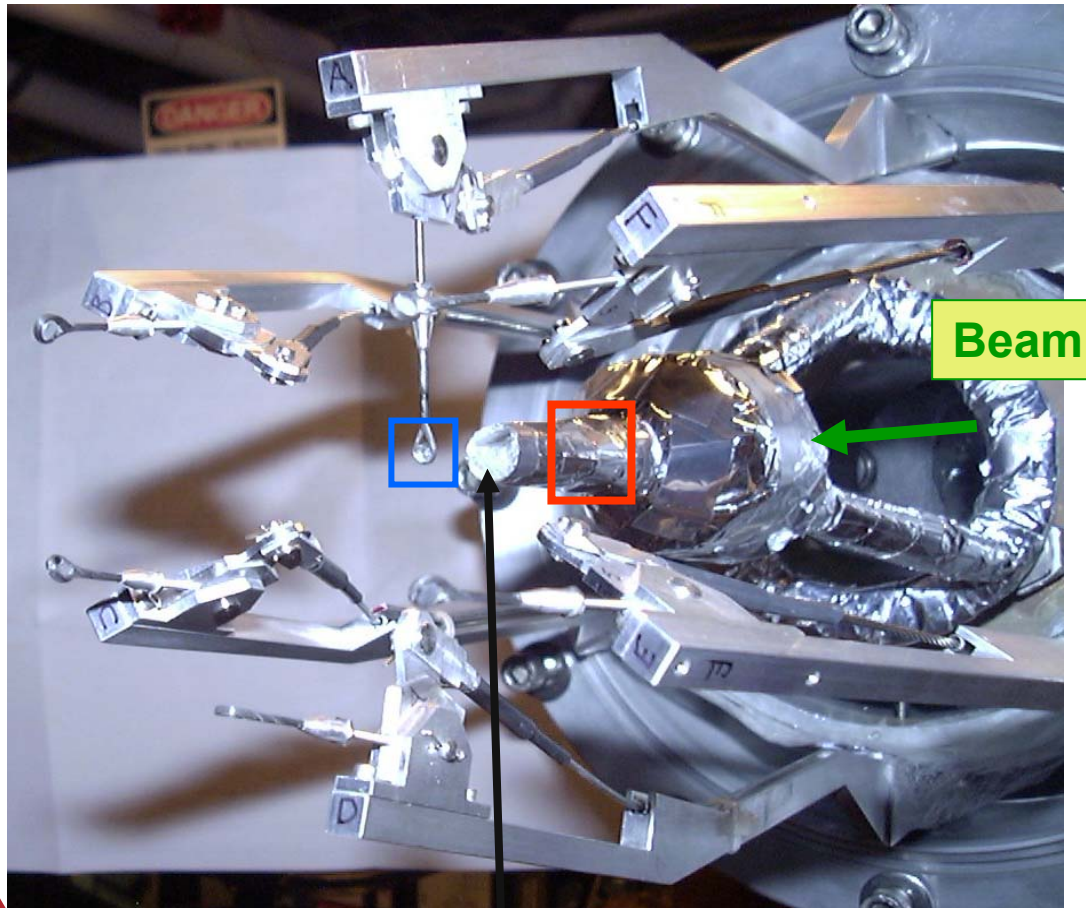


HERMES Nitrogen data : $T_A = P_0 + P_2 Q^2$

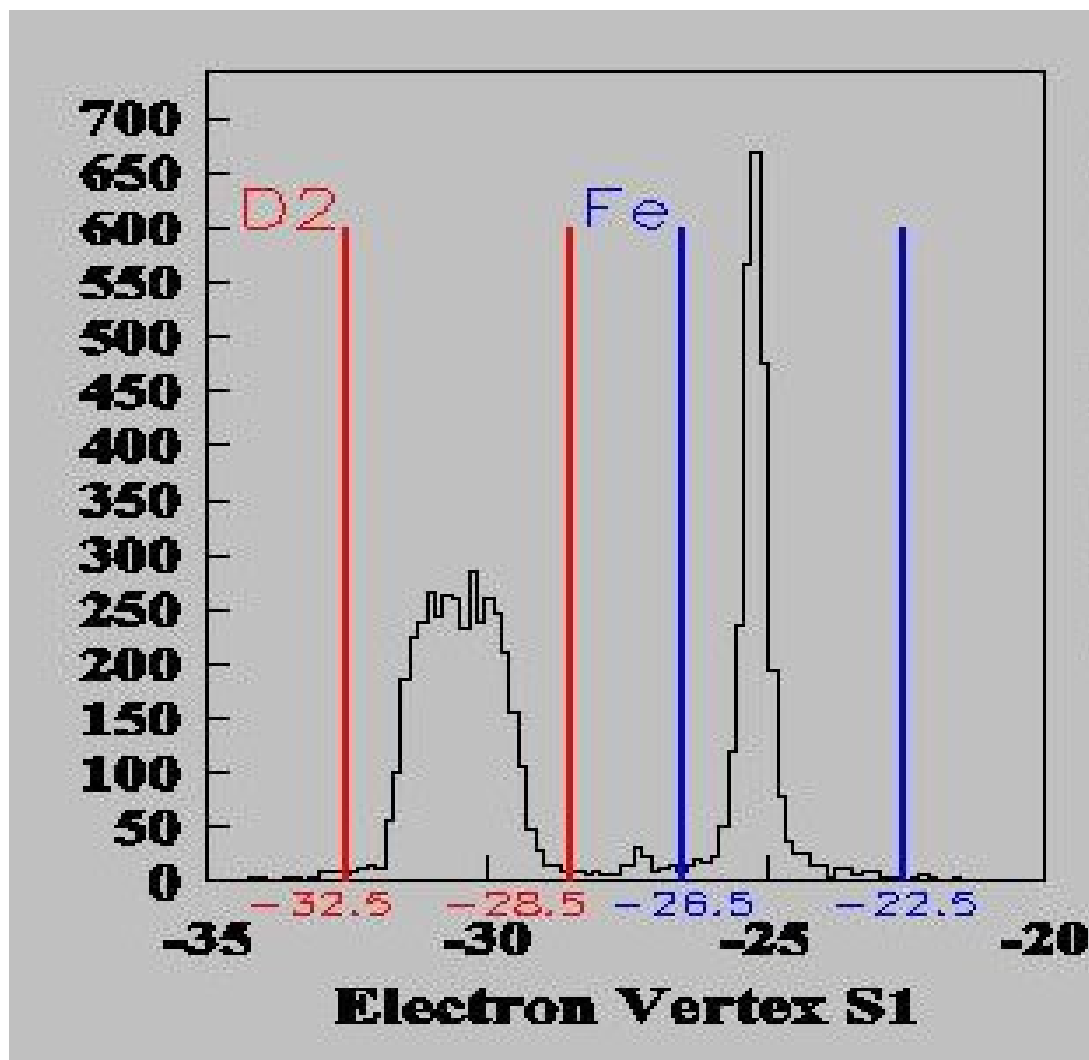
$P_2 = (0.097 \pm 0.048_{\text{stat}} \pm 0.008_{\text{syst}}) \text{ GeV}^{-2}$

JLab-CLAS E02-110
projected uncertainties

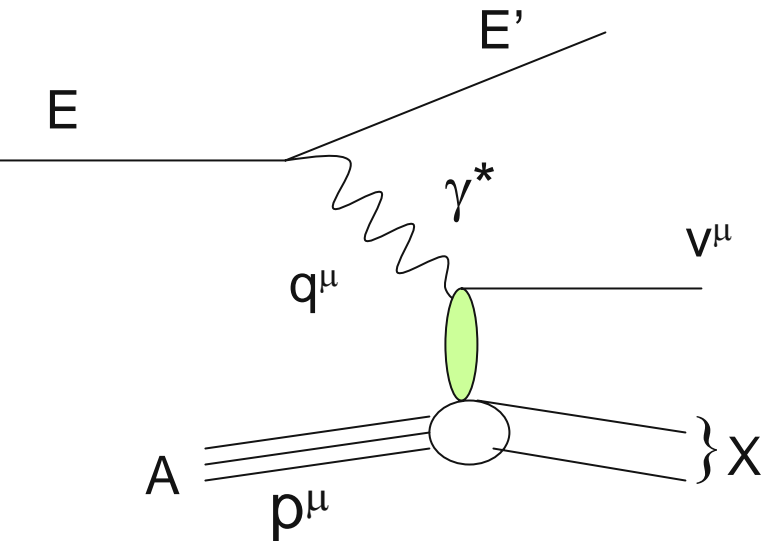
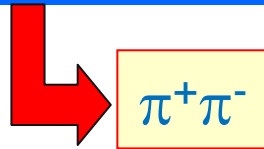
Targets



Vertex cut



$$e + \text{Fe} \rightarrow e' + \rho^0 + X$$



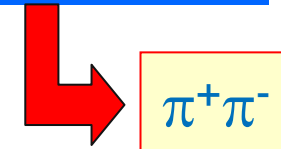
$$\nu = E - E'$$

$$Q^2 = -(q^\mu)^2 \cong 4 E E' \sin^2(\theta/2)$$

$$t = (q^\mu - v^\mu)^2$$

$$W^2 = (p^\mu + q^\mu)^2 = -Q^2 + M_p^2 + 2M_p \nu$$

$$e + D \rightarrow e' + \rho^0 + X$$



- $W \geq 2 \text{ GeV}$

\Rightarrow avoid resonance region

- $-t \leq 0.45 \text{ GeV}^2$

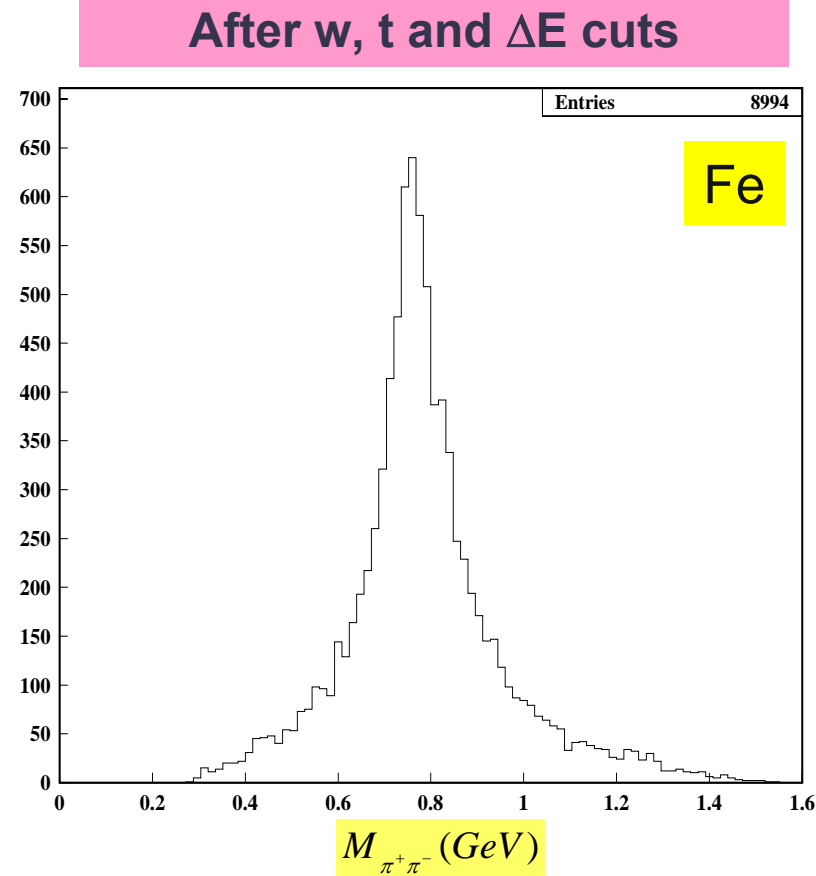
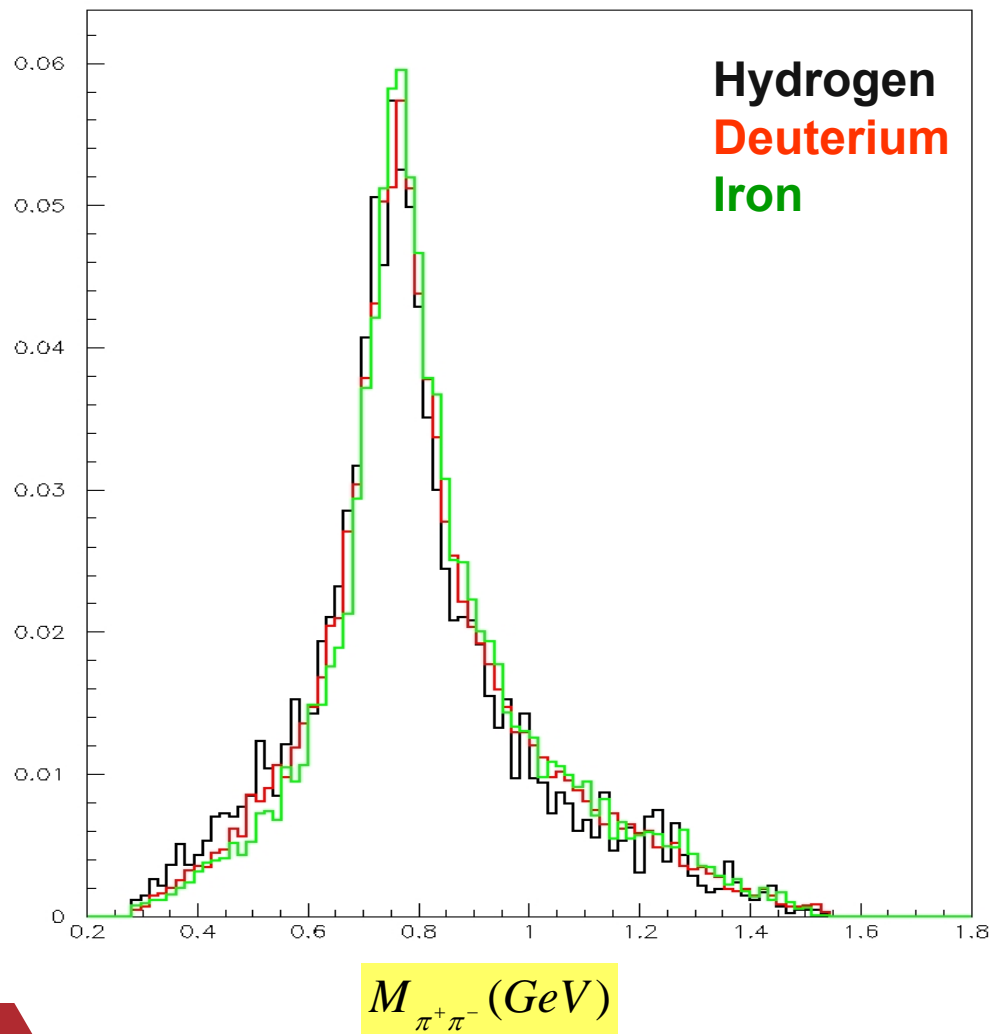
\Rightarrow select diffractive process

- $|\Delta E| \leq 0.1 \text{ GeV}$

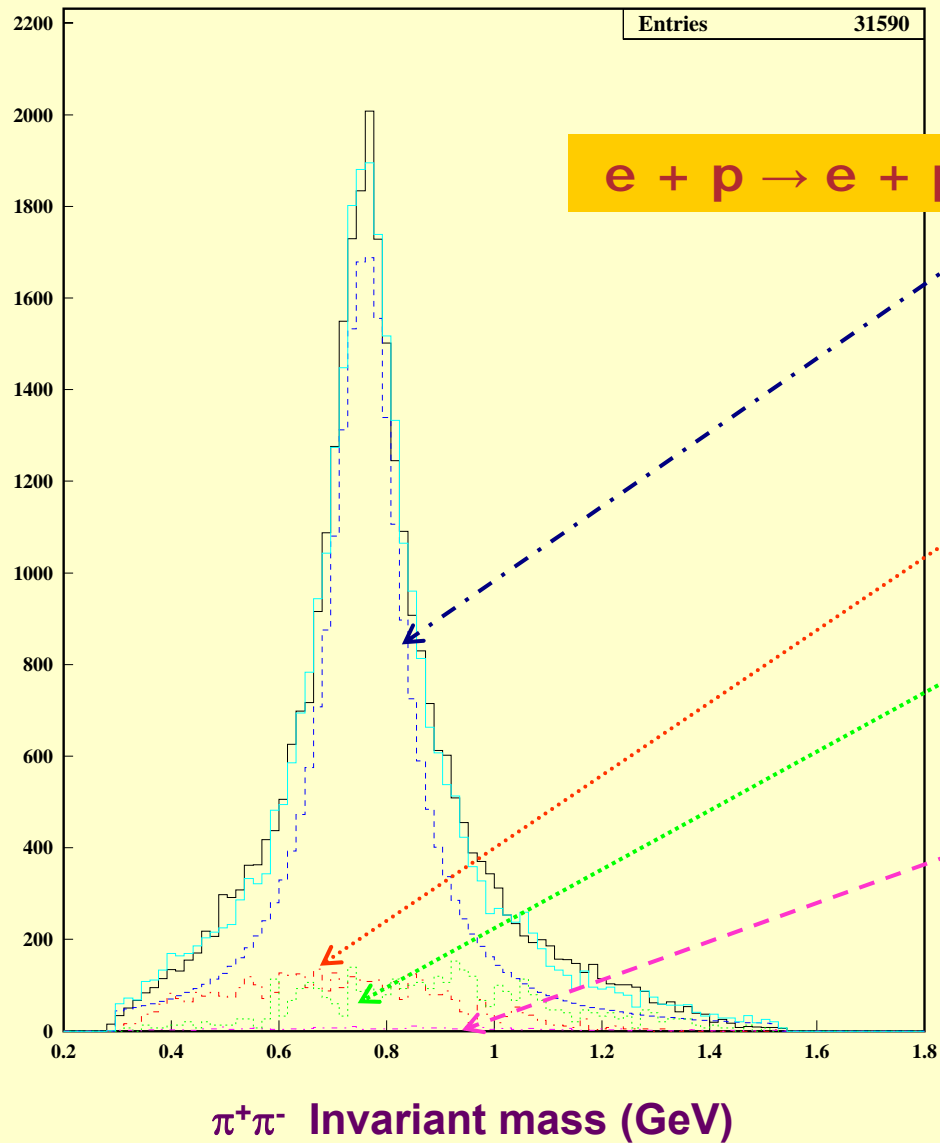
\Rightarrow select exclusive channel

- $\Delta E = \nu - E_\rho + t/2M_p$ is the missing energy from $\pi^+\pi^-$ pair due to the creation of any additional final state particles

Two pions invariant mass



D2: $w > 2$, $t \geq -0.45$ and $|\Delta E| \leq 0.1$



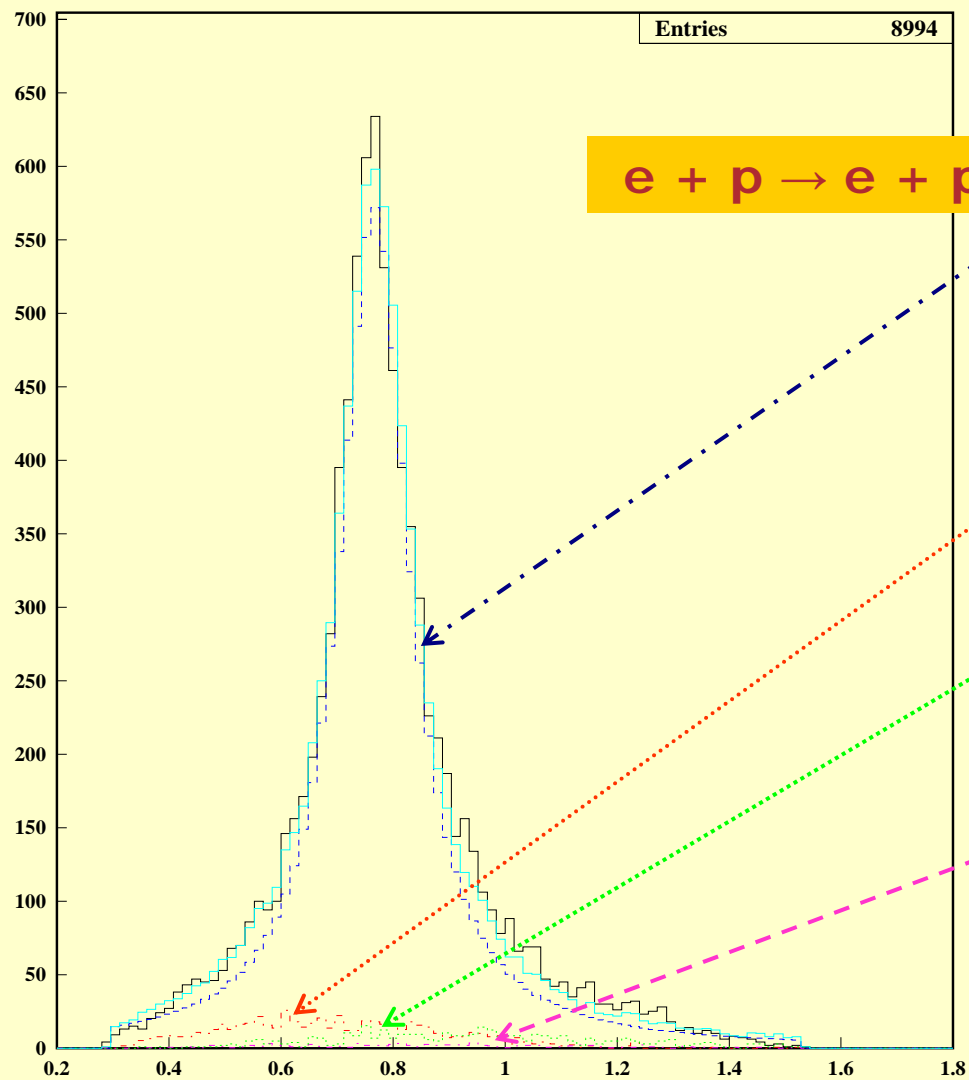
$e + p \rightarrow e + p + \rho^0$: Simple Breit Wigner

$e + p \rightarrow e + \Delta^{++} + \pi^-$

$e + p \rightarrow e + \Delta^0 + \pi^+$

$e + p \rightarrow e + p + \pi^+ + \pi^-$

Fe: $w > 2$, $t \geq -0.45$ and $|\Delta E| \leq 0.1$



$e + p \rightarrow e + p + \rho^0$: Simple Breit Wigner

$e + p \rightarrow e + \Delta^{++} + \pi^-$

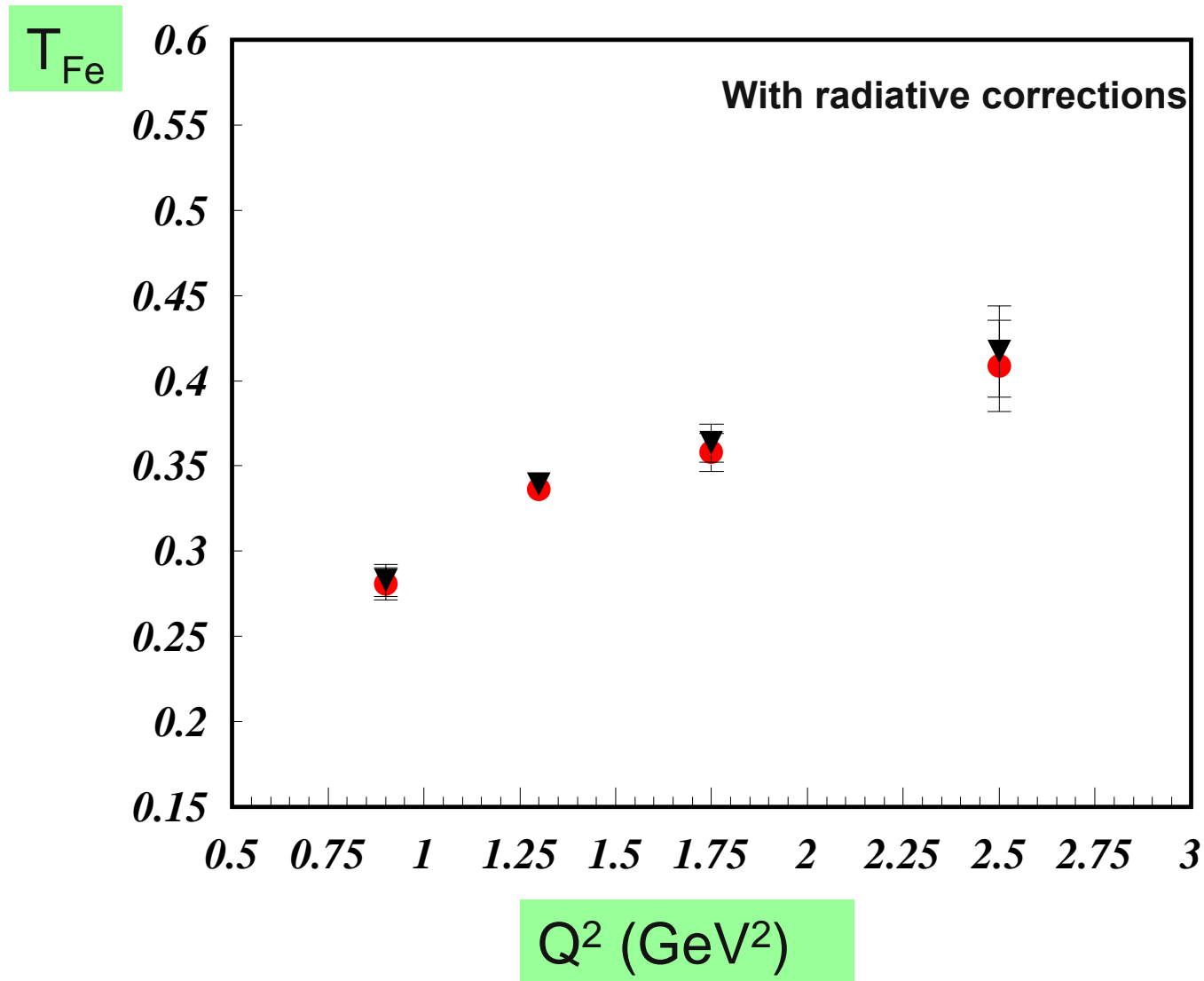
$e + p \rightarrow e + \Delta^0 + \pi^+$

$e + p \rightarrow e + p + \pi^+ + \pi^-$

$\pi^+\pi^-$ Invariant mass (GeV)

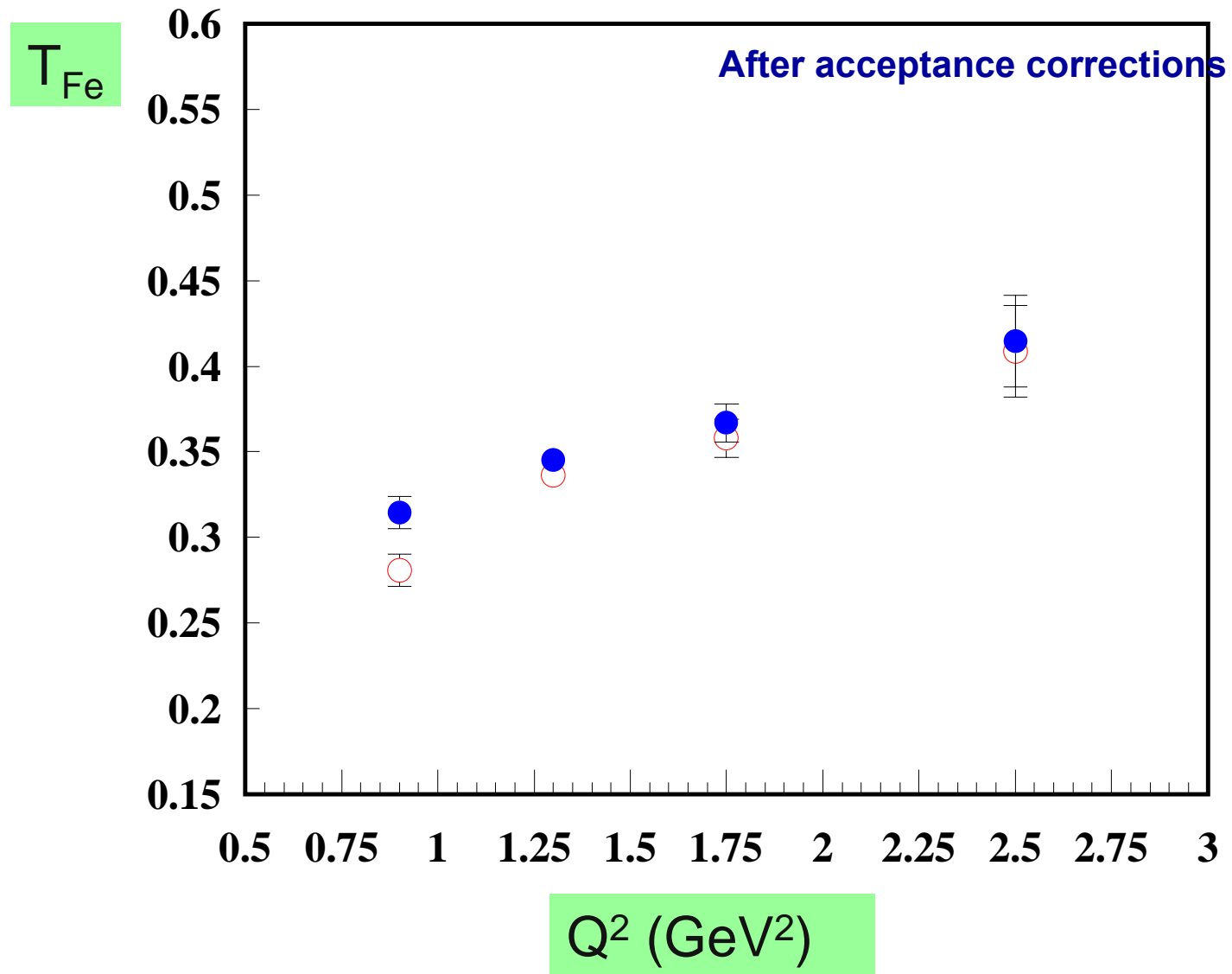
Radiative corrections

T_A vs. Q^2 corrected with radiative effect



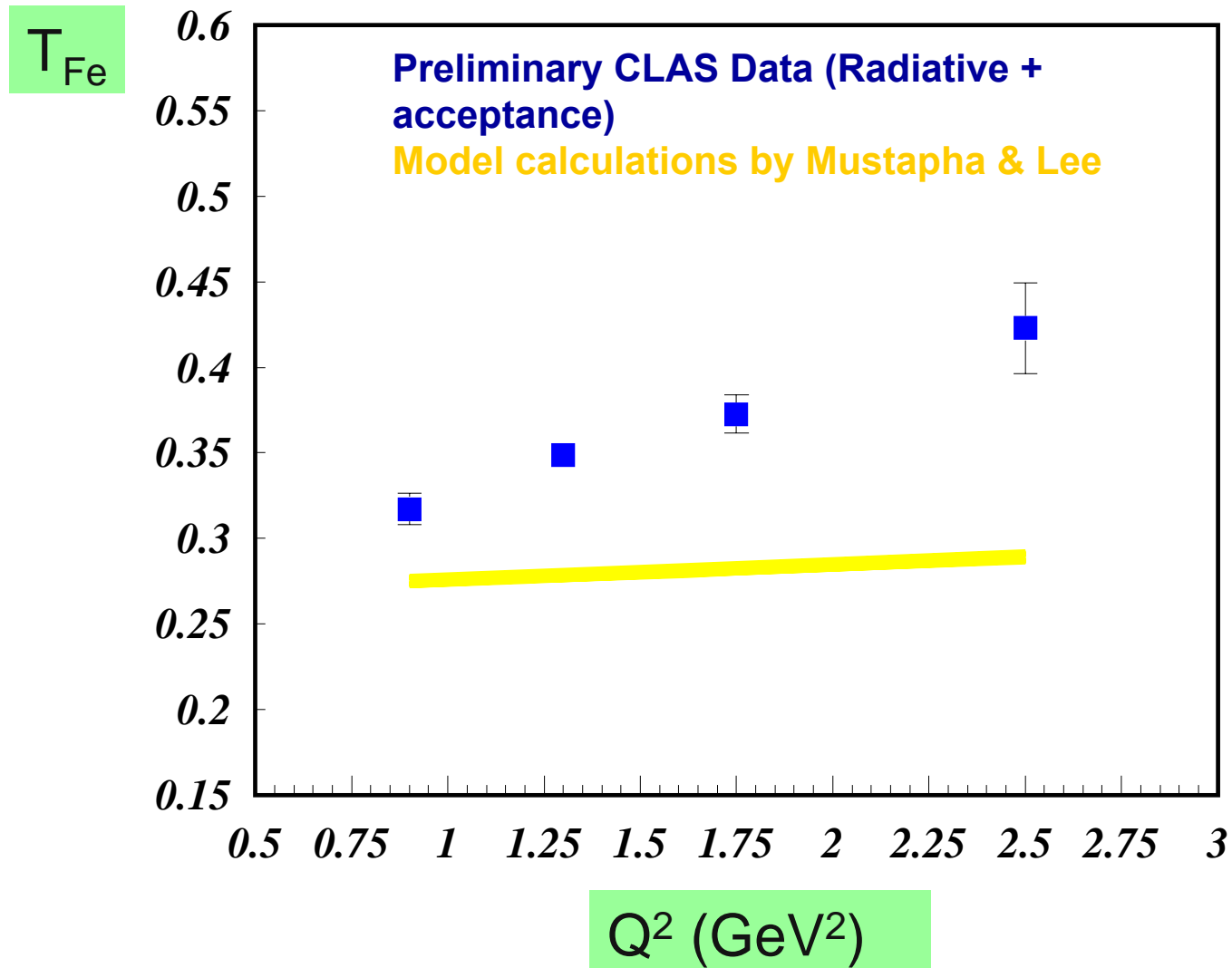
Acceptance correction

T_{Fe} vs Q^2 corrected with acceptance effect

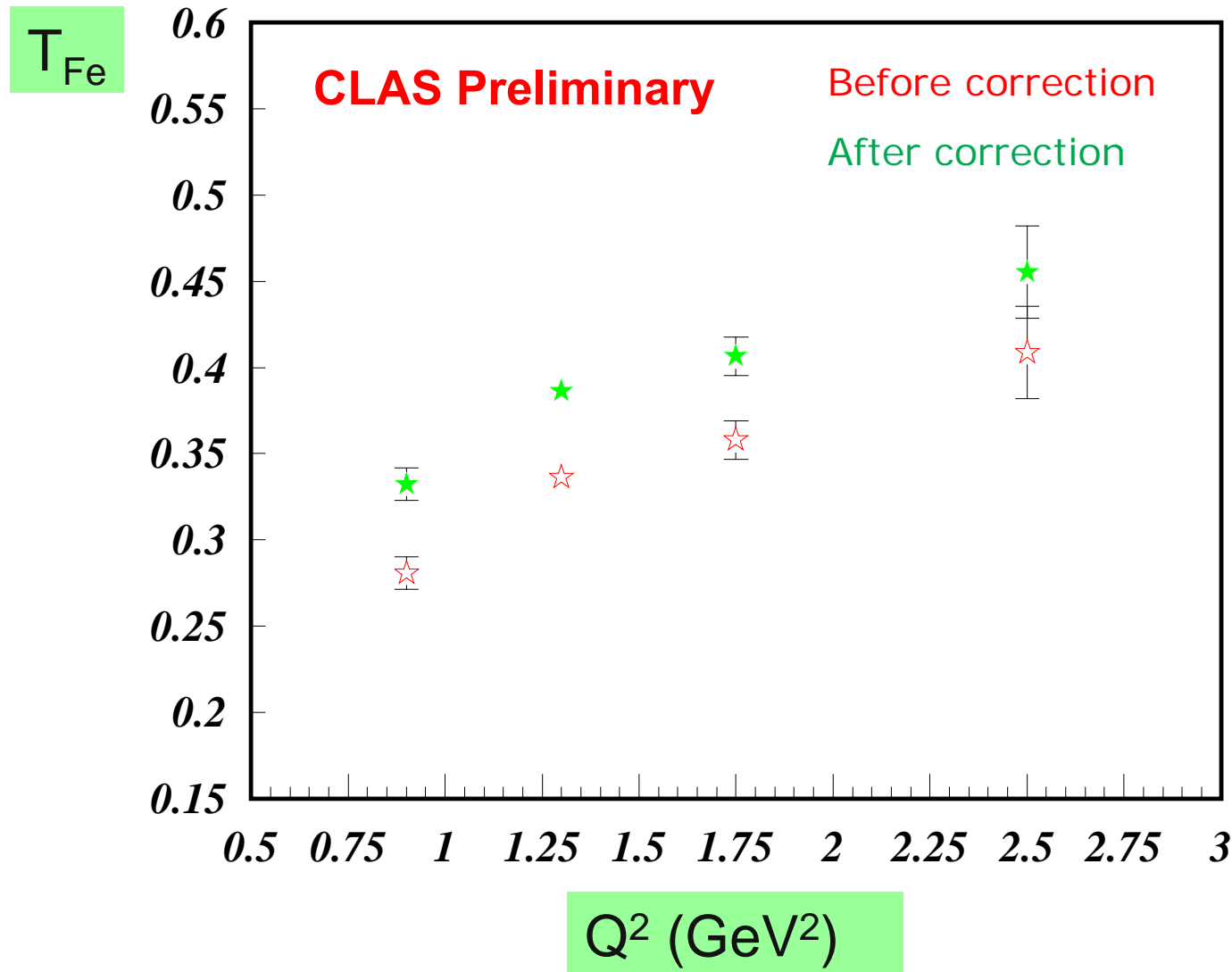


Pion absorption (Continued)

T_{Fe} vs Q^2 corrected with acceptance and radiative effects



CT Signal after Pion Absorption Correction

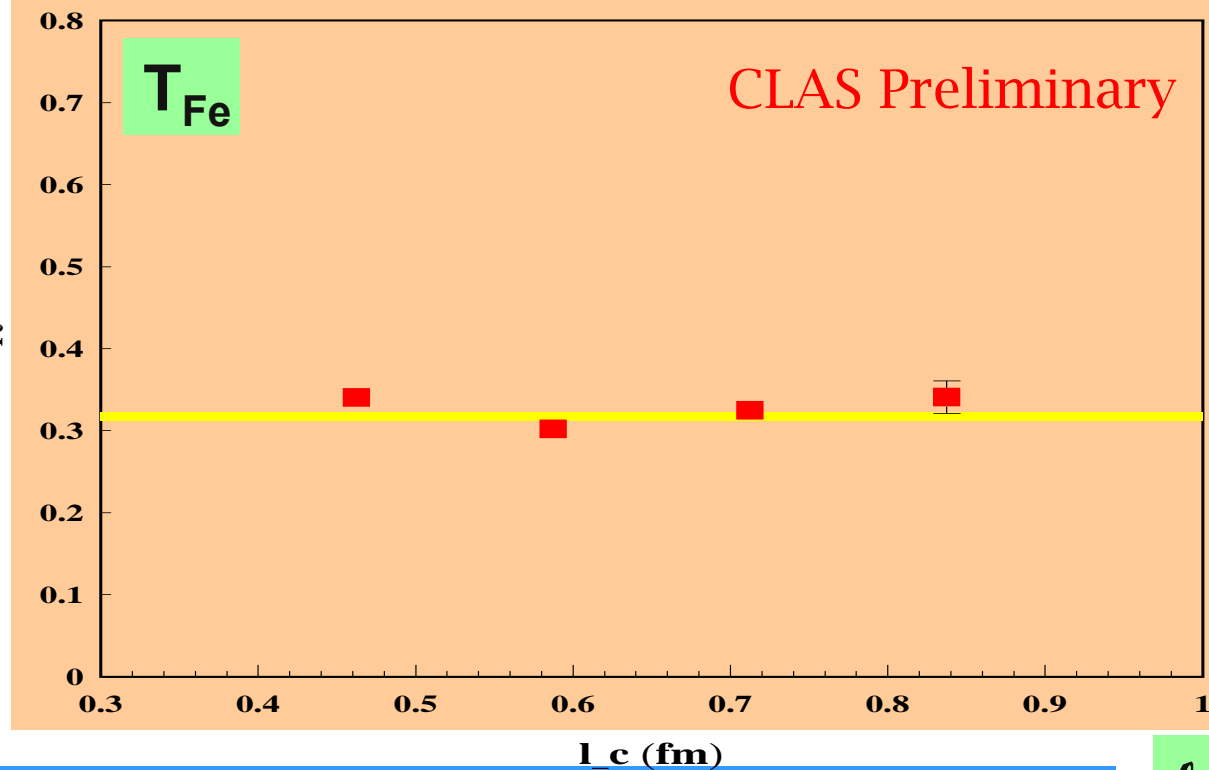


What Could mimic CT signal ?

Coherence Length

$$\ell_c = 2v/(M_v^2 + Q^2)$$

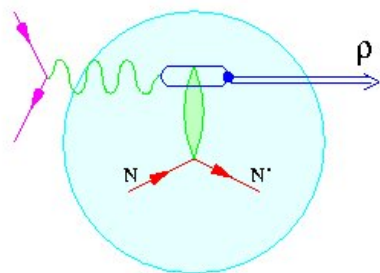
T_{Fe}^ρ



Coherence length effect (CL): Q^2 increases $\Rightarrow T_A$

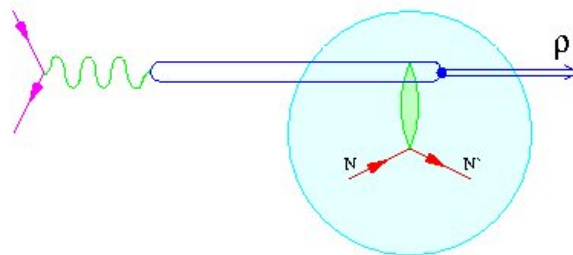
increases

Electromagnetic ISI



Small ℓ_c

Hadronic ISI



Large ℓ_c

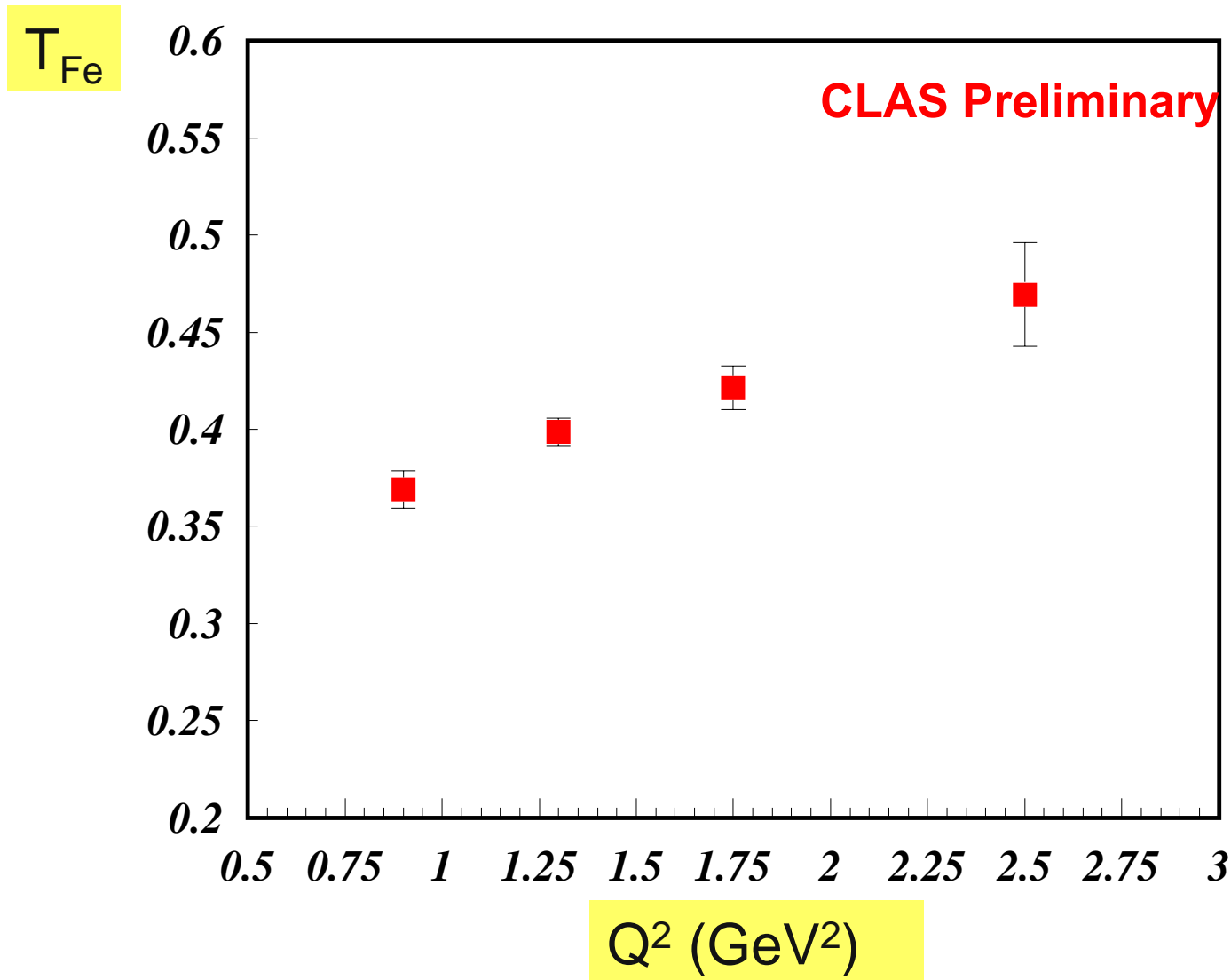
Coherence Length effect can mimic CT signal

To be safe, one should keep ℓ_c fixed and measures the Q^2 dependence of T_A

ℓ_c

Preliminary Results from CLAS EG2 data

T_A vs. Q^2 corrected with acceptance, radiative and absorption effects



Summary and Outlook

- ✿ Preliminary results from CLAS EG2 data show a strong Q^2 dependence of the nuclear transparency for Fe as predicted by the theory (B. Kopeliovich et al., Phys. Rev C 65 (2002) 035201)
- ✿ Results for 4 GeV iron and 5 GeV carbon are coming soon
- ✿ Work on systematic uncertainties is underway